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Range Of Anisotropy**

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**Abstract**

We consider local minimizers  $u : \mathbb{R}^2 \supset \Omega \rightarrow \mathbb{R}^M$  of the variational integral  $\int_{\Omega} H(\nabla u) dx$  with density  $H$  growing at least quadratically and allowing a very large scale of anisotropy. We discuss higher integrability properties of  $\nabla u$  as well as the differentiability of  $u$  in the classical sense. Moreover, a Liouville-type theorem is established.

# 1 Introduction

The aim of our note is to give a further analysis of the regularity properties such as local higher integrability of the gradient or differentiability in the classical sense of local minimizers  $u : \mathbb{R}^n \supset \Omega \rightarrow \mathbb{R}^M$  from a suitable energy space of variational integrals like

$$(1.1) \quad I[u, \Omega] = \int_{\Omega} H(\nabla u) dx,$$

provided  $\Omega$  is a domain in  $\mathbb{R}^2$  and  $H$  is a density allowing a wide range of anisotropy, but growing at least quadratically. Let us agree for the moment to the following convention:  $H$  is called an anisotropic energy density, if

$$(1.2) \quad \lambda(|Z|)|Y|^2 \leq D^2H(Z)(Y, Y) \leq \Lambda(|Z|)|Y|^2$$

holds for all  $Y, Z \in \mathbb{R}^{nM}$  with functions  $\lambda, \Lambda : [0, \infty) \rightarrow [0, \infty)$ , which can not be chosen in such a way that

$$c_1 \leq \Lambda/\lambda \leq c_2$$

is true with positive constants  $c_1$  and  $c_2$ . Starting with the pioneering work of Marcellini [Ma1–4] various authors exhibited sufficient conditions on  $\lambda, \Lambda$  and  $H$  implying

- I. full interior regularity, if  $M = 1$  (scalar case)  
or if  $H = H(|\nabla u|)$  (dependence on the modulus) and
- II. interior partial regularity in case  $n \geq 3$  together with  $M \geq 2$ .

For I. we again refer to Marcellini’s papers and his recent collaboration with Papi [MP]. We also mention the contributions of Choe [Ch], Fusco and Sbordone [FS] and of Mingione and Siepe [MS] as well as the references quoted by these authors. Further contributions to I. are given in [ABF] and [Fu]. In connection with II. the reader should consult for example the papers of Acerbi and Fusco [AF], Cupini, Guidorzi and Mascolo [CGM], Esposito, Leonetti and Mingione [ELM1,2] and of Passarelli Di Napoli and Siepe [PS] together with the references cited by these authors. We further refer to [BF1,2]. One very popular hypothesis for proving the results stated in I. and II. is the assumption

of anisotropic  $(p, q)$ -growth, which means that the functions  $\lambda$  and  $\Lambda$  occurring in (1.2) behave like

$$(1.3) \quad \lambda(t) \approx t^{p-2}, \quad \Lambda(t) \approx t^{q-2}$$

with exponents  $1 < p \leq q < \infty$ . If  $p$  and  $q$  are too far apart, then Giaquinta's counterexample [Gi] shows that even in the scalar case singular minimizers can occur. On the other hand, if we assume (1.3), then I. and II. are true, provided we additionally impose a bound of the form

$$(1.4) \quad q < c(n)p,$$

where  $c(n)$  can be chosen rather large for low dimensions  $n$ , but  $c(n) \rightarrow 1$  as  $n \rightarrow \infty$ . Let us remark that it is possible to drop (1.4) in the case  $M = 1$  and to weaken this condition for  $M \geq 2$ , if  $H$  is of splitting type and if we restrict ourselves to locally bounded local minimizers (cf. [BFZ], [BF3-5]).

Up to now the dimension  $n$  was arbitrary, but the experience in regularity theory gives rise to the hope, that in case  $n = 2$  better results for vectorial minimizers can be obtained, which means that in principle we expect the same behaviour as in the scalar case. In fact, if (1.2) holds with  $\lambda, \Lambda$  from (1.3), then in [BF6] we could show interior  $C^1$ -regularity under the assumption

$$(1.5) \quad q < 2p,$$

and for the splitting case this is even true for any exponents  $2 \leq p \leq q < \infty$  without the limitation (1.5), we refer to [BF3]. Further details are presented in [BF7]. In this note we want to discuss the twodimensional case for densities  $H : \mathbb{R}^{2M} \rightarrow [0, \infty)$  of class  $C^2$  assuming that  $p$  is equal to 2, but imposing no upper bound on  $D^2H$ . It turns out that under these weak assumptions minimizers already belong to the Sobolev class  $W_{2,\text{loc}}^2(\Omega; \mathbb{R}^M)$ . If in addition we have (1.2) and (1.3) with  $p = 2$  and  $q$  arbitrary large, then there is an open set  $\Omega_0 \subset \Omega$  (depending on the minimizer  $u$ ) such that  $u \in C^1(\Omega_0; \mathbb{R}^M)$  and  $\mathcal{H} - \dim(\Omega - \Omega_0) = 0$ . Let us now give a precise formulation of our assumptions imposed on  $H$ : without loss of generality let  $H(0) = 0$  and  $DH(0) = 0$ . Moreover it should hold:

$$(A1) \quad \exists a > 0 : |\xi| |DH(\xi)| \leq aH(\xi), \quad \xi \in \mathbb{R}^{2M};$$

$$(A2) \quad \lambda_0 := \inf_{\xi \in \mathbb{R}^{2M}} \frac{|DH(\xi)|}{|\xi|} > 0;$$

$$(A3) \quad \exists A > 0 : \frac{|DH(\xi)|}{|\xi|} \leq A \inf_{|\sigma|=1} D^2H(\xi)(\sigma, \sigma), \quad \xi \in \mathbb{R}^{2M}.$$

**REMARK 1.1.** From (A2) and (A3) it follows that we have the first inequality in (1.2) with  $\lambda(t) \equiv \lambda_0/A$ , which corresponds to the case “ $p = 2$ ”. In particular there exists a positive constant  $c$  such that

$$(1.6) \quad H(\xi) \geq c|\xi|^2, \quad \xi \in \mathbb{R}^{2M}.$$

On the other hand, hypothesis (A1) provides an upper bound for  $H(\xi)$ , more precisely we have

$$(1.7) \quad H(\xi) \leq C(|\xi|^a + 1), \quad \xi \in \mathbb{R}^{2M},$$

for a suitable constant  $C$  and with exponent  $a$  from (A1). Combining (1.6) and (1.7) we see that  $a \geq 2$ , and both inequalities show that the density  $H$  itself is of anisotropic  $(2, a)$ -growth. In order to prove (1.7) we choose  $\xi \in \mathbb{R}^{2M}$  such that  $|\xi| \geq 1$  and let  $H_0 := \sup_{|\sigma|=1} H(\sigma)$ . Then it holds

$$\begin{aligned} \ln H(\xi) &\leq \ln H(\xi) - \ln H\left(\frac{\xi}{|\xi|}\right) + \ln H_0 \\ &= \int_1^{|\xi|} \frac{d}{dt} \ln H\left(t \frac{\xi}{|\xi|}\right) dt + \ln H_0 \\ &= \int_1^{|\xi|} \frac{1}{H\left(t \frac{\xi}{|\xi|}\right)} DH\left(t \frac{\xi}{|\xi|}\right) : \frac{\xi}{|\xi|} dt + \ln H_0 \\ &\stackrel{(A1)}{\leq} \int_1^{|\xi|} \frac{a}{t} dt + \ln H_0 = a \ln |\xi| + \ln H_0, \quad \text{i.e.} \\ H(\xi) &\leq \exp(a \ln |\xi| + \ln H_0) = H_0 |\xi|^a, \end{aligned}$$

and (1.7) is established.

**REMARK 1.2.** With respect to the growth condition (1.6) a natural space for local minimizers  $u$  of the functional  $I$  defined in (1.1) is the class

$$\mathcal{C} := \{v \in W_{2,\text{loc}}^1(\Omega; \mathbb{R}^M) : I[v, \Omega'] < \infty \quad \forall \Omega' \Subset \Omega\},$$

where here and in what follows  $W_{p,\text{loc}}^k(\dots)$  is the Sobolev space of functions as introduced for example in Adam's book [Ad]. By definition  $u \in \mathcal{C}$  is a local  $I$ -minimizer iff  $I[u, \Omega'] \leq I[v, \Omega']$  holds for all subdomains  $\Omega' \Subset \Omega$  and any function  $v \in \mathcal{C}$  such that  $\text{spt}(u - v) \subset \Omega'$ .

**REMARK 1.3.** Obviously  $H$  is a (strictly) convex function, thus in addition to (A1) we have the upper bound  $H(\xi) \leq |DH(\xi)||\xi|$  for all  $\xi \in \mathbb{R}^{2M}$ .

Our results are as follows:

**THEOREM 1.1.** Let  $H$  satisfy (A1-3) and consider a local  $I$ -minimizer  $u \in \mathcal{C}$ . Then it holds:

- a)  $\int_{\Omega'} D^2 H(\nabla u)(\partial_\alpha \nabla u, \partial_\alpha \nabla u) dx \leq c(\Omega') < \infty$  for any subdomain  $\Omega' \Subset \Omega$  and for  $\alpha = 1, 2$ .  
 In particular  $u$  is of class  $W_{2,\text{loc}}^2(\Omega; \mathbb{R}^M)$ , and this is also true in the case  $\Omega \subset \mathbb{R}^n$  with  $n \geq 3$ .

b)  $\int_{\Omega_1} H(\nabla u)^2 dx \leq c(\Omega_1, \Omega_2) (\int_{\Omega_2} H(\nabla u) dx)^2$  for arbitrary subdomains  $\Omega_1 \Subset \Omega_2 \Subset \Omega$ .

**THEOREM 1.2.** *In addition to (A1-3) suppose that*

$$(A4) \quad D^2 H(\xi)(\sigma, \sigma) \leq \Lambda_0 (1 + |\xi|^2)^{\frac{q-2}{2}} |\sigma|^2, \xi, \sigma \in \mathbb{R}^{2M};$$

is true for some  $\Lambda_0 > 0$  and an exponent  $q \in [2, \infty)$ . Let  $u \in \mathcal{C}$  denote a local minimizer. Then there is an open subset  $\Omega_0$  of  $\Omega$ , whose complement is of zero Hausdorff-dimension, and  $u \in C^{1,\mu}(\Omega_0; \mathbb{R}^M)$  for any  $\mu < 1$ .

**REMARK 1.4.** *We conjecture that even under the hypotheses of Theorem 1.1 we have interior  $C^1$ -regularity of local minimizers.*

From the proof of Theorem 1.1 we will deduce the following Liouville-type result:

**THEOREM 1.3.** *Suppose that  $H$  satisfies (A1-3) and let  $u$  denote an entire local minimizer of the functional  $I$ , i.e.  $u$  is a local minimizer on the whole plane. Then, if  $\int_{\mathbb{R}^2} H(\nabla u) dx$  is finite,  $u$  is a constant function. If the same situation is considered on  $\Omega = \mathbb{R}^n$  with  $n \geq 3$ , then  $u$  is affine. Moreover, if the condition of the finiteness of the energy is replaced by the requirement that the entire local minimizer  $u$  is a bounded function, then in the 2D-case  $u$  again must be constant.*

Our paper is organized as follows: in Section 2 we present some examples of densities satisfying (A1-4). In Section 3 we will prove Theorem 1.1 and 1.3, and Section 4 is devoted to the proof of Theorem 1.2.

## 2 Some examples

We start with a construction borrowed from [BF7]. Let

$$(2.1) \quad h(t) := \int_0^t s g(s) ds, \quad t \geq 0,$$

with

$$(2.2) \quad g(t) := 1 + \int_0^t \Theta(s) ds, \quad t \geq 0,$$

where the continuous function  $\Theta : [0, \infty) \rightarrow [0, \infty)$  is given as follows: suppose that we have fixed an arbitrary large exponent  $\alpha$  and a sequence  $\{a_i\}$  such that  $0 < a_i < a_{i+1}$ ,



$\lim_{i \rightarrow \infty} a_i = \infty$ . Next we choose small positive numbers  $\varepsilon_i$  with the properties  $I_i := (a_i - \varepsilon_i, a_i + \varepsilon_i) \subset [0, \infty)$ ,  $I_i \cap I_j = \emptyset$  for  $i \neq j$  and

$$(2.3) \quad \sum_{i=1}^{\infty} \varepsilon_i a_i^{\alpha-1} < \infty.$$

Then we let

$$\Theta(t) := \begin{cases} 0 & \text{on } [0, \infty) - \bigcup_{i=1}^{\infty} I_i, \\ \text{affine linear on } (a_i - \varepsilon_i, a_i) & \text{and on} \\ (a_i, a_i + \varepsilon_i) & \text{with value } a_i^{\alpha-1} \text{ at } t = a_i, \quad i \in \mathbb{N}. \end{cases}$$

**Lemma 2.1.** *The function  $h$  defined in (2.1) and (2.2) is of class  $C^2([0, \infty))$ , strictly increasing and convex together with  $h''(0) = 1$  and  $\lim_{t \searrow 0} \frac{h(t)}{t} = 0$ . With suitable positive constants  $c_\ell$  it holds for all  $t \geq 0$*

- a)  $h(2t) \leq c_1 h(t)$ ;
- b)  $\frac{h'(t)}{t} \leq h''(t)$ , and the function  $\frac{h'(t)}{t}$  is increasing;
- c)  $h'(t)t \leq c_2 h(t)$ ;
- d)  $h''(t) \leq c_3(1 + t^2)^{\frac{\alpha}{2}} \frac{h'(t)}{t}$  ;
- e)  $c_4 t^2 \leq h(t) \leq c_5 t^2$ .

Moreover it is not possible to replace  $\alpha$  in d) by a smaller number.

**REMARK 2.1.** *With  $h$  from above let us introduce the energy density  $H_0 : \mathbb{R}^{2M} \rightarrow [0, \infty)$ ,  $H_0(Z) := h(|Z|)$ . According to e)  $H_0$  is of quadratic growth, and from b) and d) using  $\frac{h'(t)}{t} \geq 1$  we infer for all  $Y, Z \in \mathbb{R}^{2M}$*

$$|Y|^2 \leq D^2 H_0(Z)(Y, Y) \leq c_3(1 + |Z|^2)^{\frac{\alpha}{2}} g(|Z|)|Y|^2.$$

Since  $g$  is a bounded function (see below), this estimate shows that  $H_0$  is  $(p, q)$ -elliptic in the sense of (1.3) with  $p := 2$  and  $q := 2 + \alpha$ .

**Proof of Lemma 2.1:** We first observe that

$$\frac{h'(t)}{t} = g(t) \leq 1 + \int_0^\infty g'(s) ds = 1 + \int_0^\infty \Theta(s) ds = 1 + \sum_{i=1}^{\infty} \varepsilon_i a_i^{\alpha-1},$$

hence by (2.3) we find a number  $g_\infty \in (0, \infty)$  such that

$$(2.4) \quad 1 = g(0) \leq g(t) \leq g_\infty, \quad t \geq 0.$$

The properties of  $h$  stated in front of a) are immediate. For a) we use (2.4), since this inequality implies  $g(2t) \leq g_\infty g(t)$ , hence  $h(2t) = \int_0^{2t} sg(s) ds = 4 \int_0^t sg(2s) ds \leq 4g_\infty \int_0^t sg(s) ds = 4g_\infty h(t)$ .

b) is elementary:  $h''(t) = \frac{d}{dt}(tg(t)) = g(t) + tg'(t) \geq g(t) = \frac{h'(t)}{t}$ .

c) is a consequence of a): we have  $h(2t) = \int_0^{2t} h'(s) ds \geq \int_t^{2t} h'(s) ds \geq th'(t)$ , hence  $th'(t) \leq c_1 h(t)$ .

The validity of d) is equivalent to the existence of a constant  $c_6$  such that

$$(2.5) \quad tg'(t) \leq c_6 t^\alpha g(t)$$

holds for all large  $t$ . The left-hand side of (2.5) equals  $t\Theta(t)$ , and according to (2.4) the right-hand side of (2.5) behaves like  $t^\alpha$ , thus our claim is immediate by the definition of  $\Theta$  and we also see that  $\alpha$  can not be replaced by a smaller exponent. Finally, the validity of e) is obvious.  $\square$

Now we can state our examples:

**Lemma 2.2.** *Let  $B : \mathbb{R}^{2M} \times \mathbb{R}^{2M} \rightarrow \mathbb{R}$  denote any symmetric, strictly positive bilinear function. With  $h$  from Lemma 2.1 we let*

$$\begin{aligned} H_1(Z) &:= B(Z, Z) + h(|Z|), \\ H_2(Z) &:= B(Z, Z) + h\left(\sqrt{B(Z, Z)}\right), \quad Z \in \mathbb{R}^{2M}. \end{aligned}$$

*Then  $H_1$  and  $H_2$  satisfy (A1-4) with  $q = \alpha + 2$ .*

With Lemma 2.1 the proof of Lemma 2.2 is immediate. The reader should note that in these examples condition (1.5) is violated, provided we choose  $\alpha \geq 2$ . In this case we can not refer to the paper [BF6], however - according to Theorem 1.1 and 1.2 - we still have some regularity results for local minima.

### 3 Proof of Theorem 1.1 and 1.3

Let the assumptions of Theorem 1.1 hold and consider a local minimizer  $u \in \mathcal{C}$ . The following calculations can be made precise by replacing derivatives through difference quotients. If in addition we have (A4), then we can work alternatively with a local regularization with exponent  $q$  (see, e.g. [BF1]) having a sufficient degree of regularity. For the particular examples involving “the function  $h$ ” from Section 2 a quadratic regularization from below can be applied (compare [BF7]). Let  $\eta \in C_0^\infty(\Omega)$ . With “ $:$ ” denoting the scalar product of matrices and using “ $\otimes$ ” as symbol for the tensor product of vectors from  $\mathbb{R}^M$  we obtain from Euler’s equation valid for  $u$  (from now on we use the convention

of summation with respect to indices repeated twice)

$$\begin{aligned} 0 &= \int_{\Omega} \partial_{\gamma} [DH(\nabla u)] : \nabla(\eta^2 \partial_{\gamma} u) \, dx \\ &= \int_{\Omega} \eta^2 D^2 H(\nabla u)(\partial_{\gamma} \nabla u, \partial_{\gamma} \nabla u) \, dx - \int_{\Omega} DH(\nabla u) : \partial_{\gamma} [\nabla \eta^2 \otimes \partial_{\gamma} u] \, dx \end{aligned}$$

and in conclusion

$$\begin{aligned} (3.1) \quad & \int_{\Omega} D^2 H(\nabla u)(\partial_{\gamma} \nabla u, \partial_{\gamma} \nabla u) \eta^2 \, dx \\ &= \int_{\Omega} DH(\nabla u) : (\partial_{\gamma} \nabla \eta^2 \otimes \partial_{\gamma} u) \, dx + \int_{\Omega} DH(\nabla u) : (\nabla \eta^2 \otimes \Delta u) \, dx \\ &=: T_1 + T_2. \end{aligned}$$

From (A1) we get

$$\begin{aligned} T_1 &\leq 2 \int_{\Omega} |DH(\nabla u)| |\nabla u| \{ |\nabla \eta|^2 + |\nabla^2 \eta| \} \, dx \\ &\leq 2a \{ \|\nabla \eta\|_{L^{\infty}(\Omega)}^2 + \|\nabla^2 \eta\|_{L^{\infty}(\Omega)} \} \int_{\text{spt} \eta} H(\nabla u) \, dx \\ &=: c(\eta) \int_{\text{spt} \eta} H(\nabla u) \, dx, \end{aligned}$$

and for  $T_2$  we observe

$$\begin{aligned} T_2 &\leq 2 \int_{\Omega} \eta |DH(\nabla u)| |\nabla^2 u| |\nabla \eta| \, dx \\ &= 2 \int_{\Omega} \eta \left( \frac{|DH(\nabla u)|}{|\nabla u|} \right)^{1/2} |\nabla^2 u| |\nabla \eta| (|DH(\nabla u)| |\nabla u|)^{1/2} \, dx \\ &\leq \varepsilon \int_{\Omega} \eta^2 \frac{|DH(\nabla u)|}{|\nabla u|} |\nabla^2 u|^2 \, dx + c(\varepsilon) \int_{\Omega} |\nabla \eta|^2 |DH(\nabla u)| |\nabla u| \, dx, \end{aligned}$$

where  $\varepsilon > 0$  is arbitrary. Recalling (A3), choosing  $\varepsilon$  sufficiently small and applying (A1) one more time, we deduce from (3.1) and the estimates from above

$$(3.2) \quad \int_{\Omega} \eta^2 D^2 H(\nabla u)(\partial_{\gamma} \nabla u, \partial_{\gamma} \nabla u) \, dx \leq c(\eta) \int_{\text{spt} \eta} H(\nabla u) \, dx.$$

This proves part a) of Theorem 1.1, since by (A2) and (A3)

$$D^2 H(\nabla u)(\partial_{\gamma} \nabla u, \partial_{\gamma} \nabla u) \geq \frac{\lambda_0}{A} |\nabla^2 u|^2.$$

Moreover we note that the right-hand side of (3.2) is finite. We emphasize that up to now we did not make use of our assumption that  $\Omega$  is a domain in  $\mathbb{R}^2$ . For b) we apply

Sobolev's inequality and get

$$\begin{aligned}
\int_{\Omega} (\eta H(\nabla u))^2 dx &\leq c \left[ \int_{\Omega} |\nabla(\eta H(\nabla u))| dx \right]^2 \\
&\leq c \left[ \int_{\Omega} |\nabla \eta| H(\nabla u) dx + \int_{\Omega} \eta |DH(\nabla u)| |\nabla^2 u| dx \right]^2 \\
&\leq c(\eta) \left[ \int_{\text{spt } \eta} H(\nabla u) dx \right]^2 + c \left[ \int_{\Omega} \eta |DH(\nabla u)| |\nabla^2 u| dx \right]^2.
\end{aligned}$$

The last integral on the right-hand side is handled with the help of (A3) and Hölder's inequality:

$$\begin{aligned}
\left[ \int_{\Omega} |DH(\nabla u)| |\nabla^2 u| \eta dx \right]^2 &\leq c(\eta) \left( \int_{\text{spt } \eta} |DH(\nabla u)| |\nabla u| dx \right) \\
&\quad \cdot \int_{\text{spt } \eta} \frac{|DH(\nabla u)|}{|\nabla u|} |\nabla^2 u|^2 dx \\
&\leq c(\eta) \left( \int_{\text{spt } \eta} |DH(\nabla u)| |\nabla u| dx \right) \int_{\text{spt } \eta} D^2 H(\nabla u)(\partial_{\gamma} \nabla u, \partial_{\gamma} \nabla u) dx.
\end{aligned}$$

Applying (3.2) (with appropriate choice of  $\eta$ ) and using (A1), we arrive at our claim

$$\int_{\Omega_1} H(\nabla u)^2 dx \leq c(\Omega_1, \Omega_2) \left( \int_{\Omega_2} H(\nabla u) dx \right)^2$$

for arbitrary subdomains  $\Omega_1 \Subset \Omega_2 \Subset \Omega$ . □

For proving Theorem 1.3 we first observe that (3.2) immediately implies (for all dimensions  $n \geq 2$ )

$$(3.3) \quad \int_{B_R(0)} D^2 H(\nabla u)(\partial_{\gamma} \nabla u, \partial_{\gamma} \nabla u) dx \leq cR^{-2} \int_{B_{2R}(0)} H(\nabla u) dx$$

for any radius  $R > 0$ . Therefore, if we assume that

$$(3.4) \quad \int_{\mathbb{R}^n} H(\nabla u) dx < \infty,$$

we get from (3.3) by passing to the limit  $R \rightarrow \infty$  that  $\nabla^2 u = 0$ . At the same time it is easy to see that the estimates for the  $2D$ -case stated after (3.2) yield

$$\int_{B_R(0)} H(\nabla u)^2 dx \leq cR^{-2} \left( \int_{B_{2R}(0)} H(\nabla u) dx \right)^2$$

thus, under the hypothesis (3.4), we must have

$$\int_{\mathbb{R}^2} H(\nabla u)^2 dx = 0$$

and therefore the Jacobian matrix of the affine linear function  $u$  actually vanishes. Assume now that in place of (3.4) we have

$$(3.5) \quad L := \|u\|_{L^\infty(\mathbb{R}^2)} < \infty.$$

We first claim that (3.5) implies

$$(3.6) \quad \lim_{R \rightarrow \infty} R^{-2} \int_{B_R(0)} H(\nabla u) \, dx = 0,$$

so that  $\nabla^2 u = 0$  will follow from (3.3) and (3.6). For proving (3.6) let us fix a number  $\tau > 1$  to be specified later. For  $\xi \in \mathbb{R}^{2M}$ ,  $|\xi| \geq 1$ , it holds on account of (A1) and (1.7)

$$\begin{aligned} |DH(\xi)|^\tau &\leq a H(\xi)^\tau |\xi|^{-\tau} = a H(\xi) H(\xi)^{\tau-1} |\xi|^{-\tau} \\ &\leq c H(\xi) |\xi|^{a(\tau-1)-\tau} \leq c H(\xi), \end{aligned}$$

provided we choose  $\tau$  such that  $a(\tau-1) - \tau \leq 0$ . In this case we arrive at

$$(3.7) \quad |DH(\xi)| \leq c (H(\xi)^{1/\tau} + 1), \quad \xi \in \mathbb{R}^{2M}.$$

Next consider  $R > 0$ , let  $k \in \mathbb{N}$  and choose  $\eta \in C_0^\infty(B_{2R}(0))$ ,  $\eta = 1$  on  $B_R(0)$ ,  $0 \leq \eta \leq 1$ ,  $|\nabla \eta| \leq c/R$ .

Starting from

$$0 = \int_{B_{2R}(0)} DH(\nabla u) : \nabla (\eta^{2k} u) \, dx,$$

observing  $H(\xi) \leq \xi : DH(\xi)$  (see Remark 1.3) and using (3.5) we find

$$(3.8) \quad \int_{B_{2R}(0)} \eta^{2k} H(\nabla u) \, dx \leq c \int_{B_{2R}(0)} \eta^{2k-1} |\nabla \eta| |DH(\nabla u)| \, dx$$

with constant  $c$  depending on  $L$  and  $k$ . On the right-hand side of (3.8) we apply Young's inequality in combination with (3.7) and get for all  $\varepsilon > 0$

$$\begin{aligned} &\int_{B_{2R}(0)} \eta^{2k-1} |\nabla \eta| |DH(\nabla u)| \, dx \\ &\leq c \int_{B_{2R}(0)} \eta^{2k-1} H(\nabla u)^{1/\tau} |\nabla \eta| \, dx + c \int_{B_{2R}(0)} |\nabla \eta| \, dx \\ &\leq \varepsilon \int_{B_{2R}(0)} \eta^{(2k-1)\tau} H(\nabla u) \, dx + c(\varepsilon) \int_{B_{2R}(0)} |\nabla \eta|^{\frac{\tau}{\tau-1}} \, dx \\ &\quad + c \int_{B_{2R}(0)} |\nabla \eta| \, dx. \end{aligned}$$

Thus, if  $\varepsilon$  is small enough and if we choose  $k$  so large that  $(2k-1)\tau \geq 2k$ , then the  $\varepsilon$ -term can be absorbed into the left-hand side of (3.8). Observing that  $\tau/(\tau-1) > 1$ , we deduce at least for  $R \geq 1$

$$\int_{B_R(0)} H(\nabla u) \, dx \leq c \int_{B_{2R}(0)} |\nabla \eta| \, dx \leq cR,$$

and (3.6) follows. From the above inequality and the estimate stated after (3.4) we infer  $\int_{\mathbb{R}^2} H(\nabla u)^2 dx < \infty$ , but since  $\nabla u$  is a constant matrix, this is only possible in the case that  $\nabla u$  vanishes.  $\square$

## 4 Proof of Theorem 1.2

Suppose now that in addition to (A1-3) the hypothesis (A4) with  $q > 2$  is valid and fix a local minimizer  $u \in \mathcal{C}$ . By Theorem 1.1 and Sobolev's embedding theorem  $|\nabla u|$  belongs to the space  $\bigcap_{1 \leq s < \infty} L_{\text{loc}}^s(\Omega)$ , thus the excess function

$$E(x, r) =: \int_{B_r(x)} |\nabla u - (\nabla u)_{x,r}|^2 dy + \int_{B_r(x)} |\nabla u - (\nabla u)_{x,r}|^q dy,$$

where  $(g)_{x,r}$  denotes the mean value of a function  $g$  with respect to a disc  $B_r(x) \Subset \Omega$ , is well defined. We claim:

**Lemma 4.1.** *Fix  $L > 0$ . Then there exists a constant  $C_*(L)$  such that for every  $\tau \in (0, 1/4)$  there is an  $\varepsilon = \varepsilon(L, \tau)$  satisfying: if  $B_r(x) \Subset \Omega$  and if we have*

$$|(\nabla u)_{x,r}| \leq L, \quad E(x, r) \leq \varepsilon(L, \tau),$$

then  $E(x, \tau r) \leq C_*(L)E(x, r)\tau^2$ .

**Proof:** We argue by contradiction following the ideas of [BF1] assuming that  $L > 0$  is fixed. The constant  $C_*(L)$  will be specified below. If the lemma is wrong, then for some  $\tau$  there are discs  $B_{r_n}(x_n) \Subset \Omega$  such that

$$(4.1) \quad |(\nabla u)_{x_n, r_n}| \leq L, \quad E(x_n, r_n) =: \lambda_n^2 \rightarrow 0, \quad n \rightarrow \infty,$$

$$(4.2) \quad E(x_n, r_n \tau) > C_* \tau^2 \lambda_n^2.$$

Letting  $a_n := (u)_{x_n, r_n}$ ,  $A_n := (\nabla u)_{x_n, r_n}$  and

$$u_n(z) := \frac{1}{\lambda_n r_n} [u(x_n + r_n z) - a_n - r_n A_n z], \quad |z| < 1,$$

we obtain from (4.1) and (4.2)

$$(4.3) \quad |A_n| \leq L, \quad \int_{B_1} |\nabla u_n|^2 dz + \lambda_n^{q-2} \int_{B_1} |\nabla u_n|^q dz = 1.$$

$$(4.4) \quad \int_{B_\tau} |\nabla u_n - (\nabla u_n)_{0,\tau}|^2 dx + \lambda_n^{q-2} \int_{B_\tau} |\nabla u_n - (\nabla u_n)_{0,\tau}|^q dz > C_* \tau^2.$$

From (4.3) we get after passing to subsequences

$$(4.5) \quad \begin{aligned} u_n &\rightharpoonup \bar{u} && \text{in } W_2^1(B_1; \mathbb{R}^M), \quad A_n \rightharpoonup \bar{A}, \\ \lambda_n \nabla u_n &\rightarrow 0 && \text{in } L^2(B_1; \mathbb{R}^{2M}) \text{ and a.e.}, \\ \lambda_n^{1-2/q} \nabla u_n &\rightarrow 0 && \text{in } L^q(B_1; \mathbb{R}^{2M}). \end{aligned}$$

It is easy to verify that  $\bar{u}$  satisfies

$$\int_{B_1} D^2 H(\bar{A})(\nabla \bar{u}, \nabla \varphi) dx = 0, \quad \varphi \in C_0^\infty(B_1; \mathbb{R}^M),$$

and thereby is a smooth function for which the Campanato estimate

$$(4.6) \quad \int_{B_\tau} |\nabla \bar{u} - (\nabla \bar{u})_{0,\tau}|^2 dx \leq C^* \tau^2$$

holds with a suitable constant  $C^* = C^*(L)$ . Let us choose  $C_* := 2C^*$ . Then (4.4) and (4.6) are in contradiction, if we can improve the weak convergences from (4.5) to

$$(4.7) \quad \nabla u_n \rightarrow \nabla \bar{u} \text{ in } L_{\text{loc}}^2(B_1; \mathbb{R}^{2M}),$$

$$(4.8) \quad \lambda_n^{1-2/q} \nabla u_n \rightarrow 0 \text{ in } L_{\text{loc}}^q(B_1; \mathbb{R}^{2M}).$$

The claim (4.7) follows exactly as (4.16) i) in [BF1] by quoting Proposition 4.3 from this reference (letting  $\mu = 0$  there). For verifying (4.8) we let

$$\Psi_n := \lambda_n^{-1} [(1 + |A_n + \lambda_n \nabla u_n|^2)^{1/2} - (1 + |A_n|^2)^{1/2}]$$

and observe that from  $(Q \in \mathbb{R}^{2M}, \eta \in C_0^\infty(\Omega))$

$$0 = \int_{\Omega} \partial_\gamma [DH(\nabla u)] : \nabla(\eta^2 \partial_\gamma [u - Qx]) dx$$

we obtain the Caccioppoli inequality

$$(4.9) \quad \int_{\Omega} \eta^2 D^2 H(\nabla u)(\partial_\gamma \nabla u, \partial_\gamma \nabla u) dx \leq c \int_{\Omega} |\nabla \eta|^2 |D^2 H(\nabla u)| |\nabla u - Q|^2 dx.$$

Combining (4.9) with (A2) and (A3) we get after scaling for  $\rho \in (0, 1)$  (choosing  $Q = A_n$ )

$$(4.10) \quad \int_{B_\rho} |\nabla \Psi_n|^2 dz \leq c(\rho) \int_{B_1} |D^2 H(A_n + \lambda_n \nabla u_n)| |\nabla u_n|^2 dx.$$

Now, by (4.3), the right-hand side of (4.10) is bounded through a finite constant so that

$$(4.11) \quad \sup_n \int_{B_\rho} |\nabla \Psi_n|^2 dz \leq c(\rho) < \infty, \quad 0 < \rho < 1.$$

At the same time we have

$$|\Psi_n| \leq c |\nabla u_n|$$

and therefore in addition to (4.11) (recall (4.3))

$$(4.12) \quad \sup_n \int_{B_\rho} |\Psi_n|^2 dz \leq c(\rho) < \infty, \quad 0 < \rho < 1.$$

With (4.11) and (4.12) it is shown that  $\{\Psi_n\}$  is a bounded sequence in each space  $W_2^1(B_\rho)$ ,  $0 < \rho < 1$ , and this will imply (4.8): to this purpose we fix a number  $K \gg 1$ . Letting

$$U_n := U_n(K, \rho) := \{z \in B_\rho : \lambda_n |\nabla u_n(z)| \leq K\}$$

we obtain from  $q > 2$  and (4.7) (using the smoothness of  $\bar{u}$ )

$$\begin{aligned} \int_{U_n} \lambda_n^{q-2} |\nabla u_n|^q dz &\leq c \left\{ \lambda_n^{q-2} \int_{U_n} |\nabla u_n - \nabla \bar{u}|^q dz + \lambda_n^{q-2} \int_{U_n} |\nabla \bar{u}|^q dz \right\} \\ &\leq c \left\{ \int_{U_n} \lambda_n^{q-2} (|\nabla u_n|^{q-2} + |\nabla \bar{u}|^{q-2}) |\nabla u_n - \nabla \bar{u}|^2 dz + \int_{U_n} \lambda_n^{q-2} |\nabla \bar{u}|^q dz \right\} \rightarrow 0 \\ \text{as } n &\rightarrow \infty. \end{aligned}$$

On the other hand, for  $K$  large enough and  $z \in B_\rho - U_n$  it holds  $\Psi_n(z) \geq c |\nabla u_n(z)|$ , hence

$$\lambda_n^{q-2} \Psi_n^q(z) \geq c \lambda_n^{q-2} |\nabla u_n(z)|^q.$$

By Sobolev's embedding theorem we have

$$\sup_n \int_{B_\rho} |\Psi_n|^q dz < \infty,$$

and since we assume  $q > 2$ , it follows from the above estimates

$$\int_{B_\rho - U_n} \lambda_n^{q-2} |\nabla u_n|^q dz \rightarrow 0, \quad n \rightarrow \infty.$$

This proves (4.8) and thereby Lemma 4.1. □

Let us define the set

$$\Omega_0 := \{x \in \Omega : \sup_{r>0} |(\nabla u)_{x,r}| < \infty \text{ and } \liminf_{r \searrow 0} E(x, r) = 0\}.$$

Then, according to Lemma 4.1,  $u$  is of class  $C^{1,\mu}$  in a neighborhood of each point  $x \in \Omega_0$ , and it remains to check more precisely, which points  $x \in \Omega$  belong to the set  $\Omega_0$ . Since  $u$  is in  $W_{2,\text{loc}}^2(\Omega; \mathbb{R}^M)$ , we deduce from Poincaré's inequality

$$\begin{aligned} \int_{B_r(x)} |\nabla u - (\nabla u)_{x,r}|^2 dy &\leq cr^2 \int_{B_r(x)} |\nabla^2 u|^2 dy \\ &= c \int_{B_r(x)} |\nabla^2 u|^2 dy \rightarrow 0, \quad r \rightarrow 0, \end{aligned}$$



for all points  $x \in \Omega$ . Next let  $s := \frac{2q}{2+q}$ . The Sobolev-Poincaré estimate gives

$$\left( \int_{B_r(x)} |\nabla u - (\nabla u)_{x,r}|^q dy \right)^{1/q} \leq cr \left( \int_{B_r(x)} |\nabla^2 u|^s dy \right)^{1/s},$$

hence we find after applying Hölder's inequality

$$\begin{aligned} \int_{B_r(x)} |\nabla u - (\nabla u)_{x,r}|^q dy &\leq cr^q r^{-\frac{2}{s}q} \left( \int_{B_r(x)} |\nabla^2 u|^s dy \right)^{q/s} \\ c r^{q-\frac{2}{s}q} r^{2\frac{q}{s}-q} \left( \int_{B_r(x)} |\nabla^2 u|^2 dy \right)^{q/2} &\rightarrow 0, \quad r \rightarrow 0, \end{aligned}$$

so that  $\lim_{r \rightarrow 0} E(x, r) = 0$  for all  $x \in \Omega$ . Therefore  $\Omega_0 = \{x \in \Omega : \sup_{r>0} |(\nabla u)_{x,r}| < \infty\}$  and the complement of this set is of Hausdorff-dimension zero. This finishes the proof of Theorem 1.2.  $\square$

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