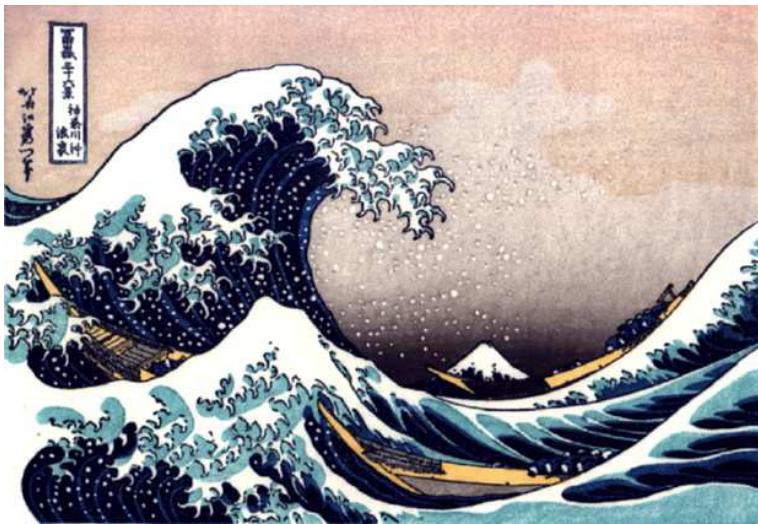
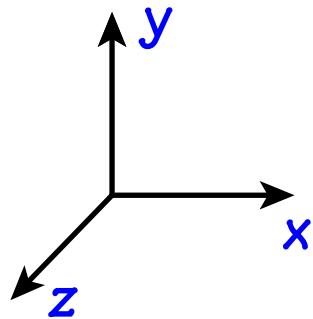


THE WATER-WAVE PROBLEM



$$y = h + \eta(x, z, t)$$



Velocity potential ϕ

Kinematic boundary condition:

$$\eta_t = \phi_y - \eta_x \phi_x - \eta_z \phi_z$$

Dynamic boundary condition:

$$\phi_t + \frac{1}{2}(\phi_x^2 + \phi_y^2 + \phi_z^2) + g\eta - \sigma \left[\frac{\eta_x}{\sqrt{1 + \eta_x^2 + \eta_z^2}} \right]_x - \sigma \left[\frac{\eta_z}{\sqrt{1 + \eta_x^2 + \eta_z^2}} \right]_z = 0$$

Difficulties:

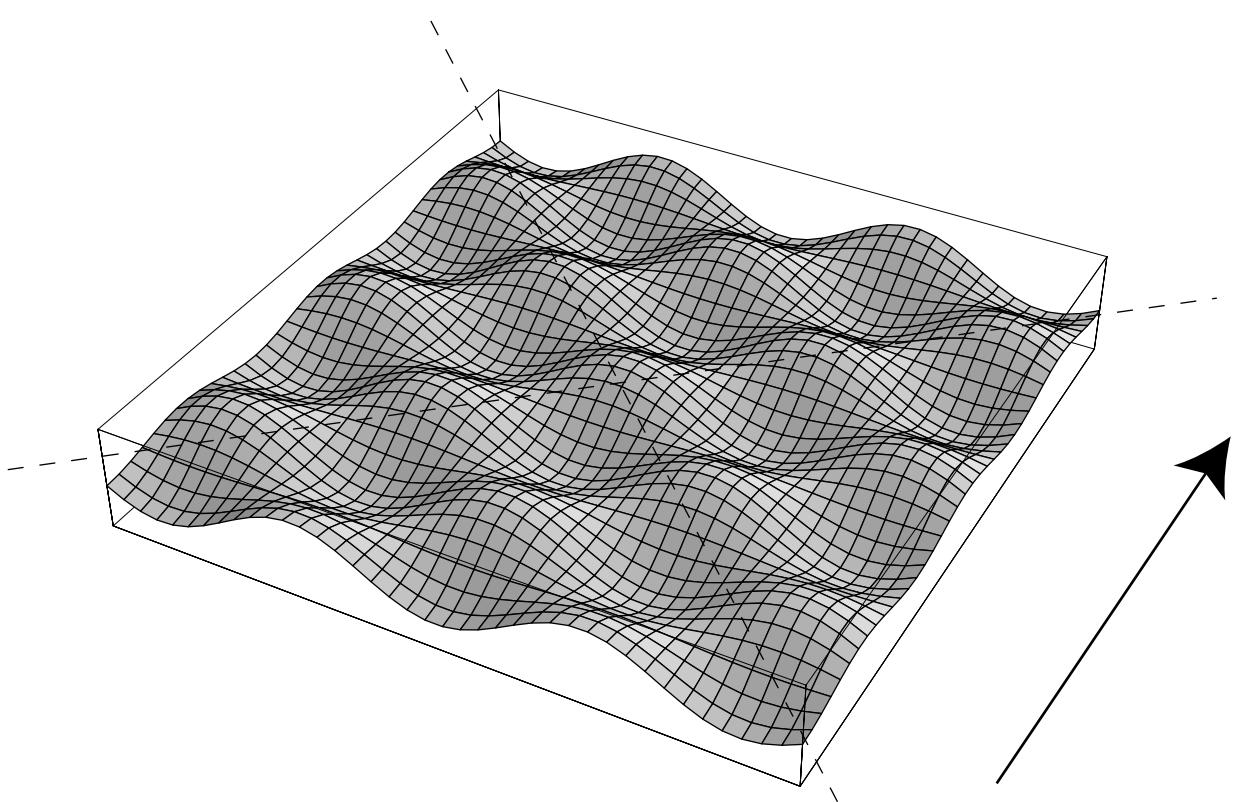
- A free boundary problem
- Nonlinear boundary conditions

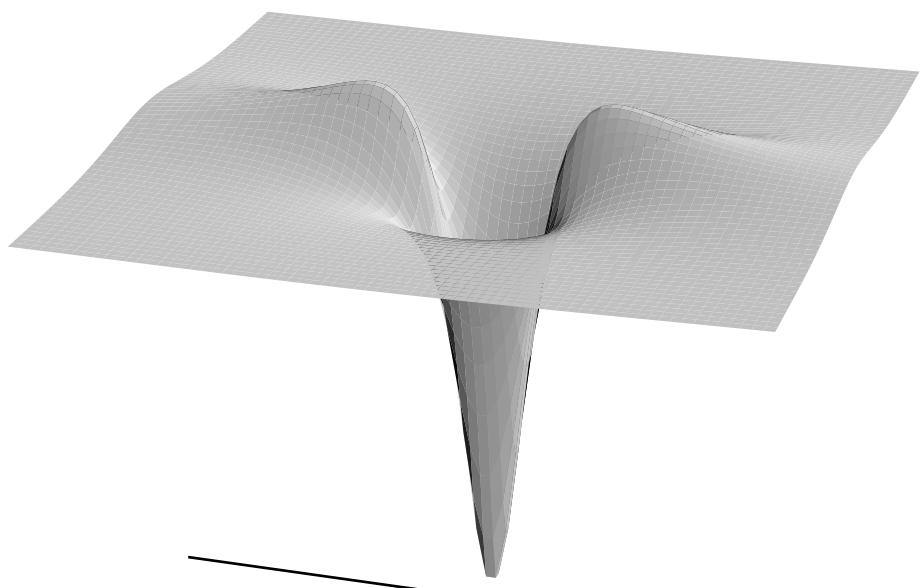
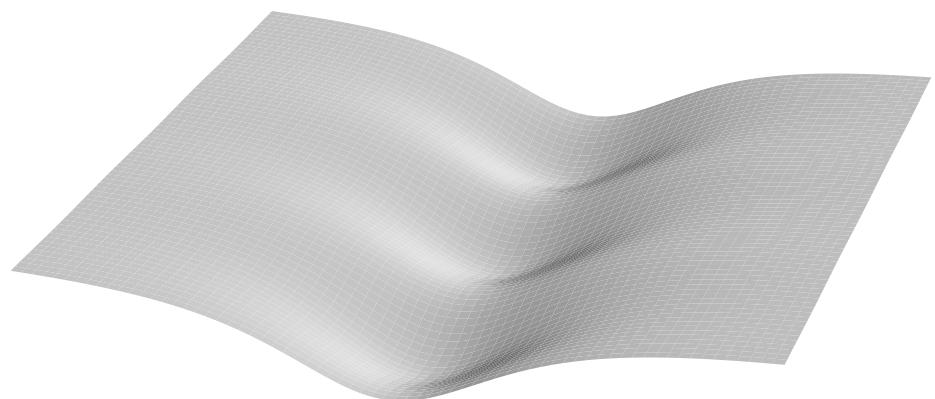
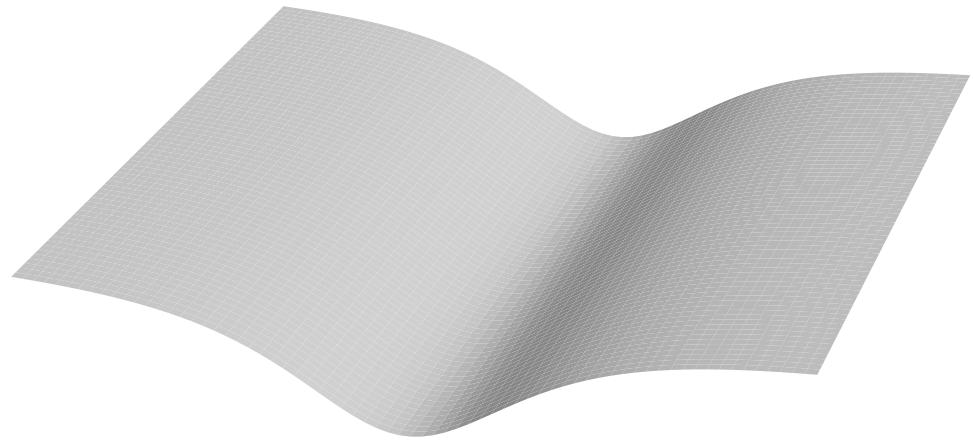
Travelling waves:

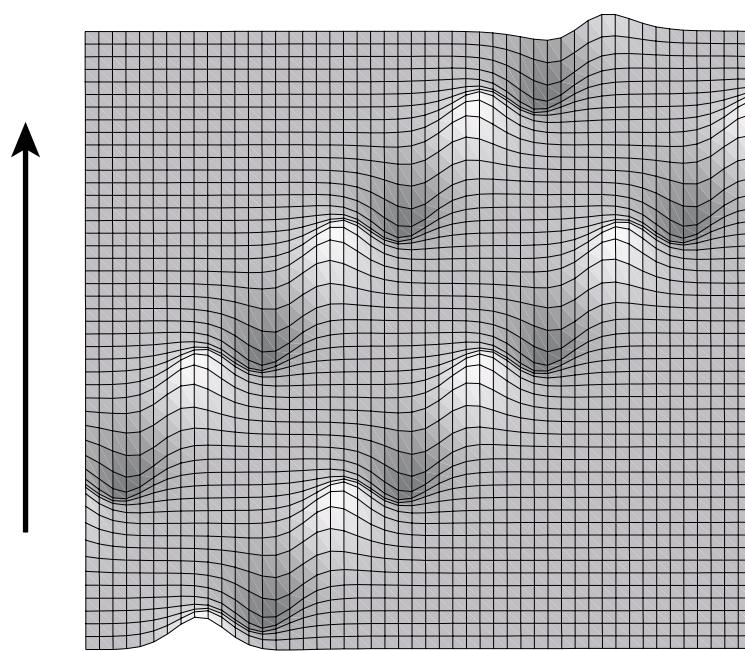
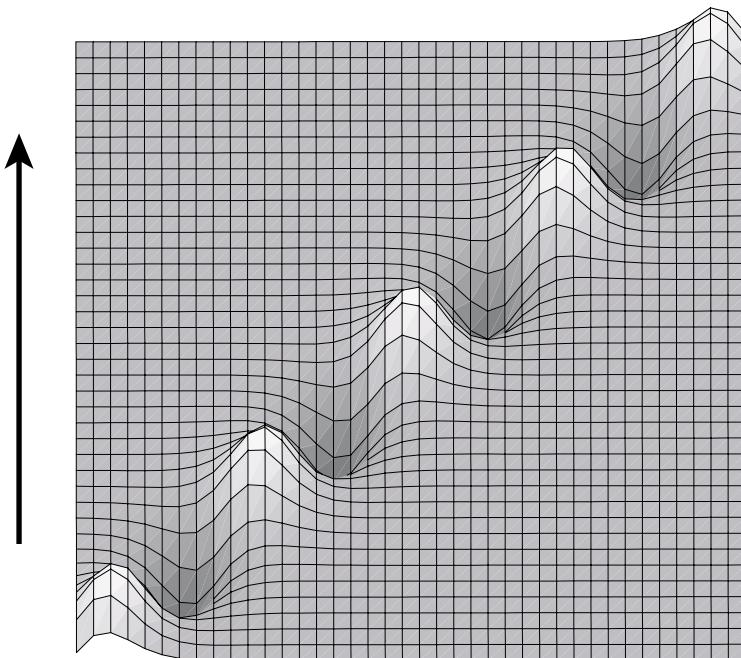
$$\eta(x, z, t) = \eta(x - ct, z), \quad \phi(x, y, z, t) = \phi(x - ct, y, z)$$

Parameters:

$$a = gh/c^2, \quad \beta = \sigma/hc^2$$







VARIATIONAL PRINCIPLE

Luke's variational principle:

$$\delta \int_{\mathbb{R}^2} \left\{ \int_0^{1+\eta} \left(-\phi_x + \frac{1}{2}(\phi_x^2 + \phi_y^2 + \phi_z^2) \right) dy + \frac{1}{2}a\eta^2 + \beta(\sqrt{1 + \eta_x^2 + \eta_z^2} - 1) \right\} dx dz = 0$$

New variables:

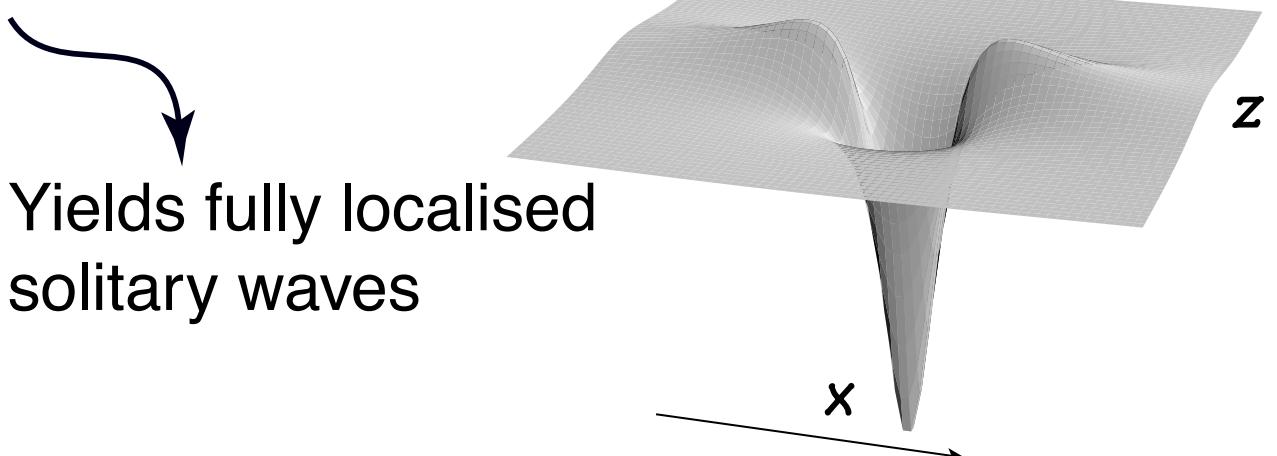
$$\tilde{y} = y/(1 + \eta(x, z)),$$

$$\phi(x, y, z) = \Phi(x, \tilde{y}, z)$$

New variational principle:

$$\delta \int_{\mathbb{R}^2} \left\{ \int_0^1 \left(\frac{1}{2} \left(\Phi_x - \frac{y\eta_x\Phi_y}{1+\eta} \right)^2 + \frac{\Phi_y^2}{2(1+\eta)^2} + \frac{1}{2} \left(\Phi_z - \frac{y\eta_z\Phi_y}{1+\eta} \right)^2 \right) (1 + \eta) dy + \eta\Phi_x|_{y=1} + \frac{1}{2}a\eta^2 + \beta[\sqrt{1 + \eta_x^2 + \eta_z^2} - 1] \right\} dx dz = 0$$

- Apply the direct methods of the calculus of variations

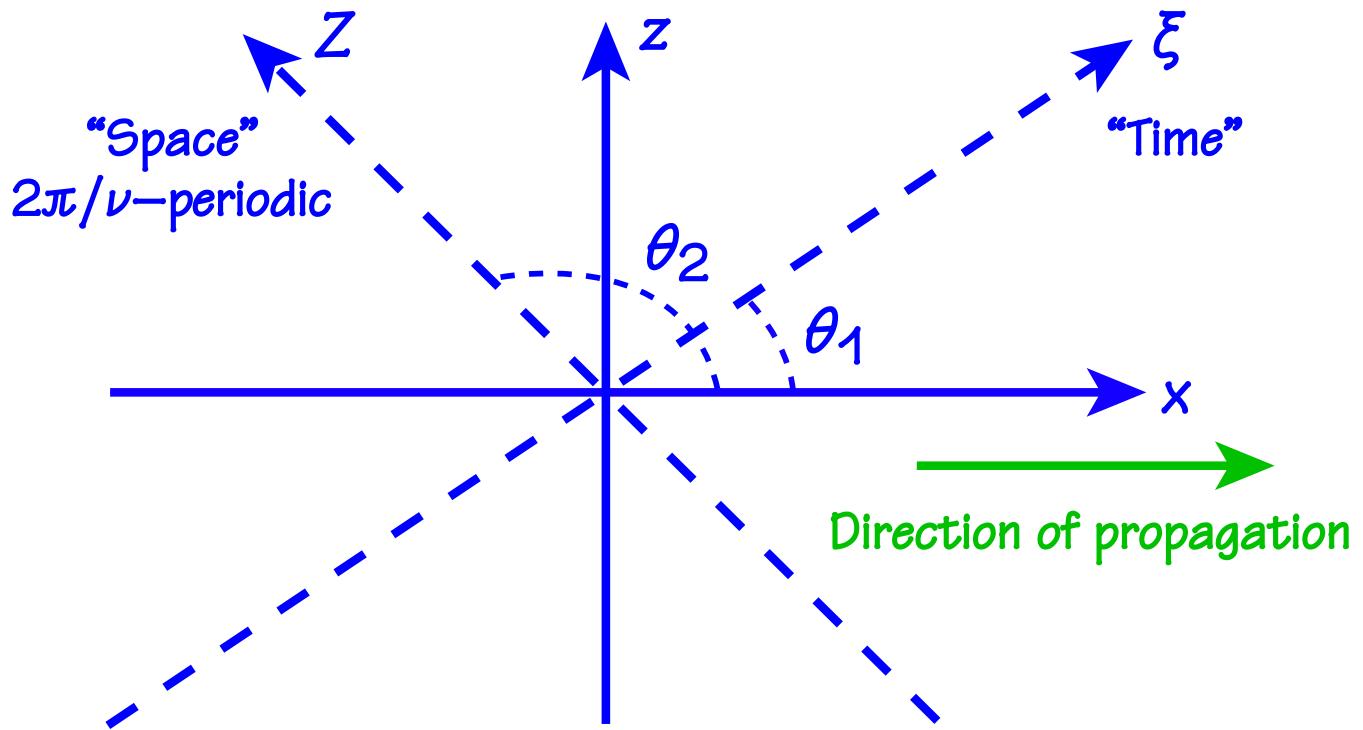


Yields fully localised
solitary waves

SPATIAL DYNAMICS

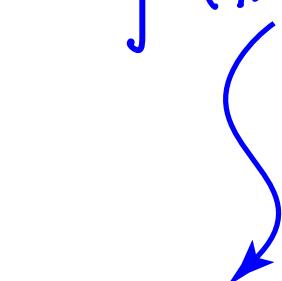
Variational principle:

$$\delta \int_{\mathbb{R}^2} \left\{ \int_0^1 \left(\frac{1}{2} \left(\Phi_x - \frac{y \eta_x \Phi_y}{1+\eta} \right)^2 + \frac{\Phi_y^2}{2(1+\eta)^2} + \frac{1}{2} \left(\Phi_z - \frac{y \eta_z \Phi_y}{1+\eta} \right)^2 \right) (1+\eta) dy + \eta \Phi_x|_{y=1} + \frac{1}{2} a \eta^2 + \beta [\sqrt{1+\eta_x^2+\eta_z^2} - 1] \right\} dx dz = 0$$



Perform the Legendre transform:

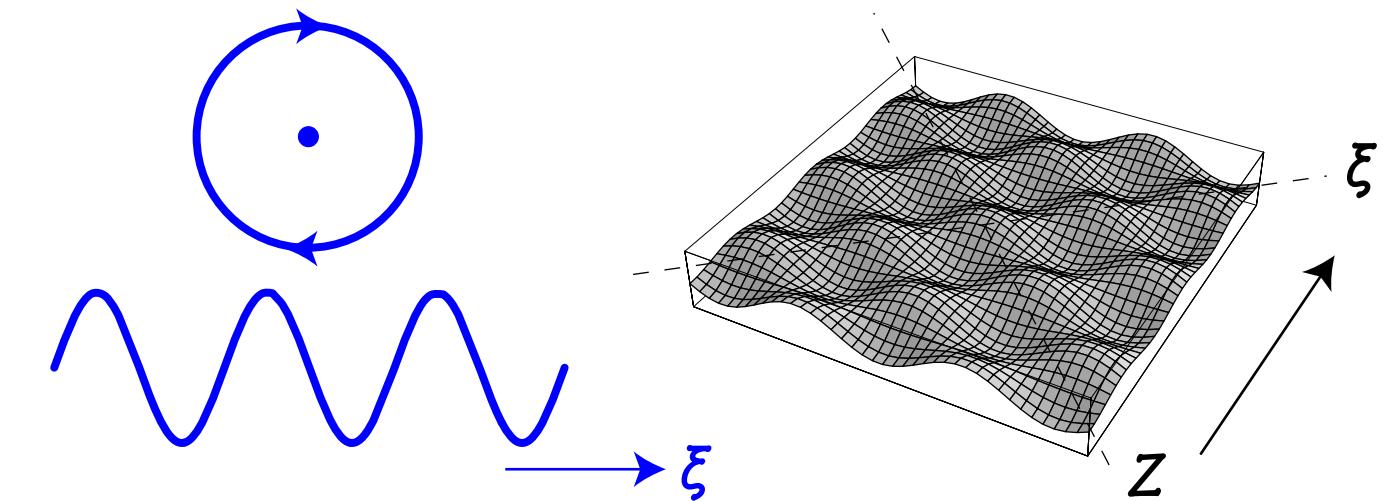
$$\delta \int f(\eta, \Phi, \eta_\xi, \Phi_\xi) d\xi = 0$$



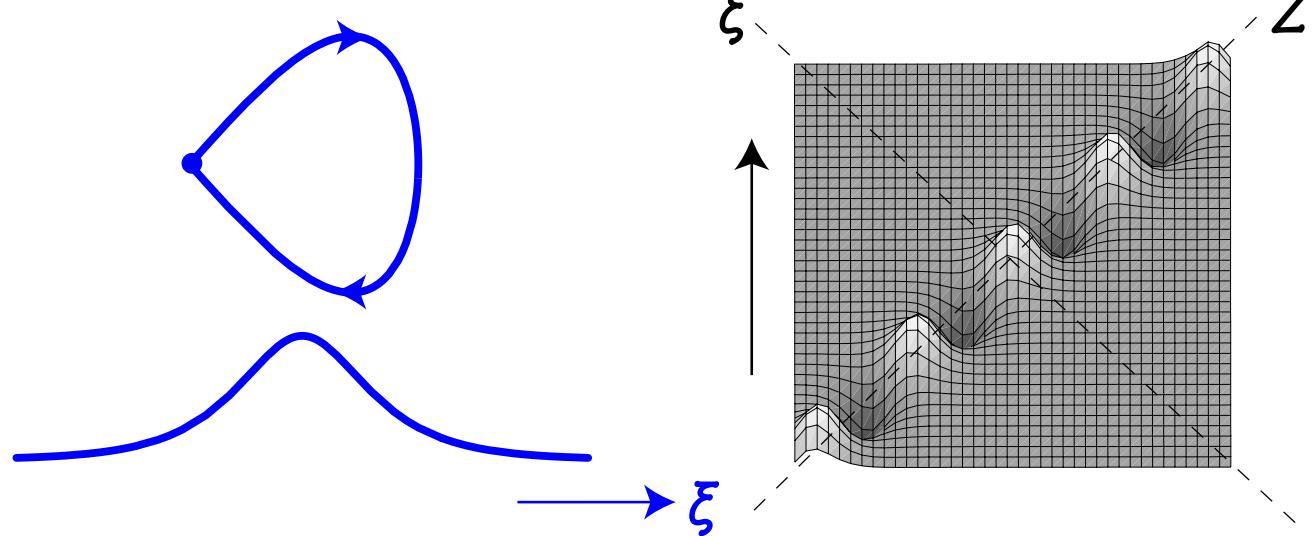
$$u_\xi = Lu + N(u), \quad u = (\eta, \Phi, \partial_{\eta_\xi} f, \partial_{\Phi_\xi} f)$$

SOLUTIONS OF THE EVOLUTIONARY EQUATION

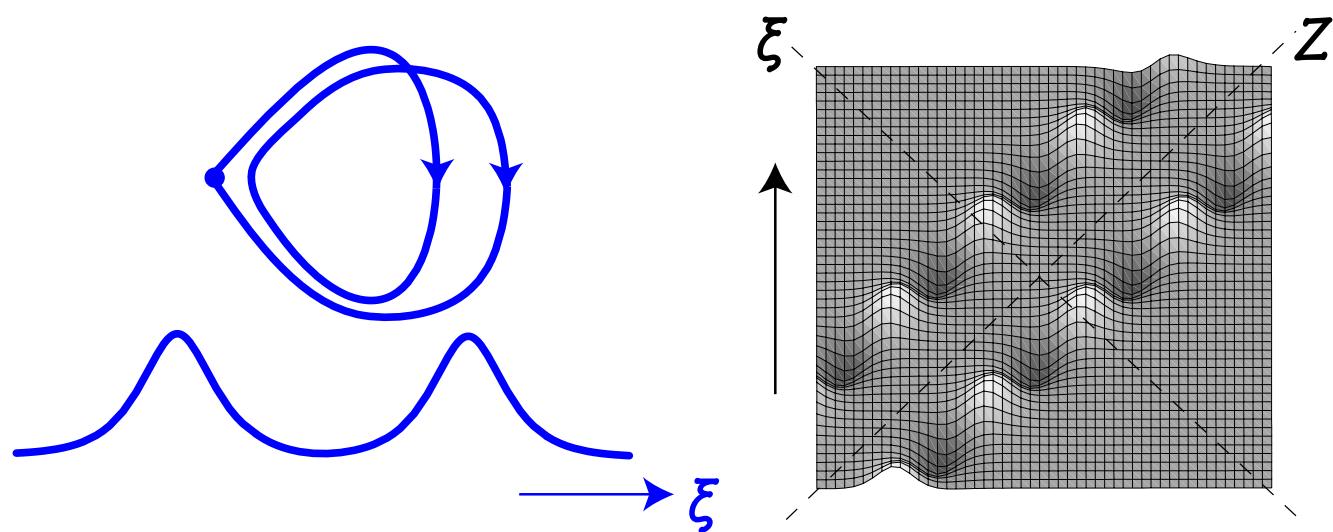
- Periodic solutions:



- Homoclinic solutions:



- Multi-pulse homoclinic solutions:



THE CENTRE MANIFOLD

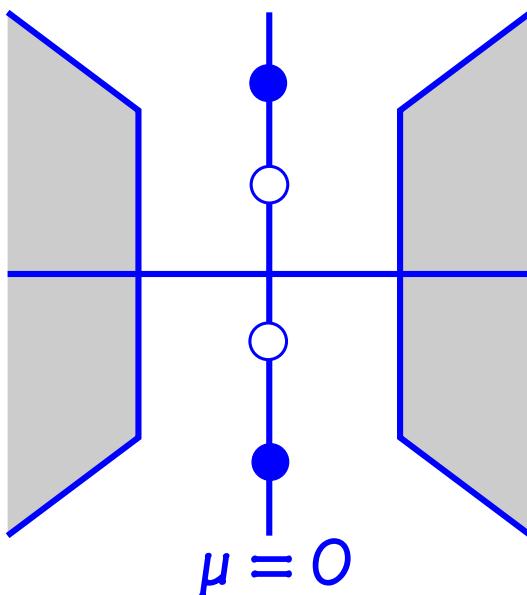
Introduce bifurcation parameters

$$(a, \beta, \nu, \theta_1, \theta_2) = (a_0, \beta_0, \nu_0, \theta_{10}, \theta_{20}) + \mu,$$

so that

$$u_\xi = L_\mu u + N_\mu(u) \quad (\star)$$

Spectrum:



- (\star) has a finite-dimensional invariant manifold M
- For small μ all small, bounded solutions to (\star) lie on M
- Reduced system:
 - describes the flow on M
 - has a Hamiltonian structure

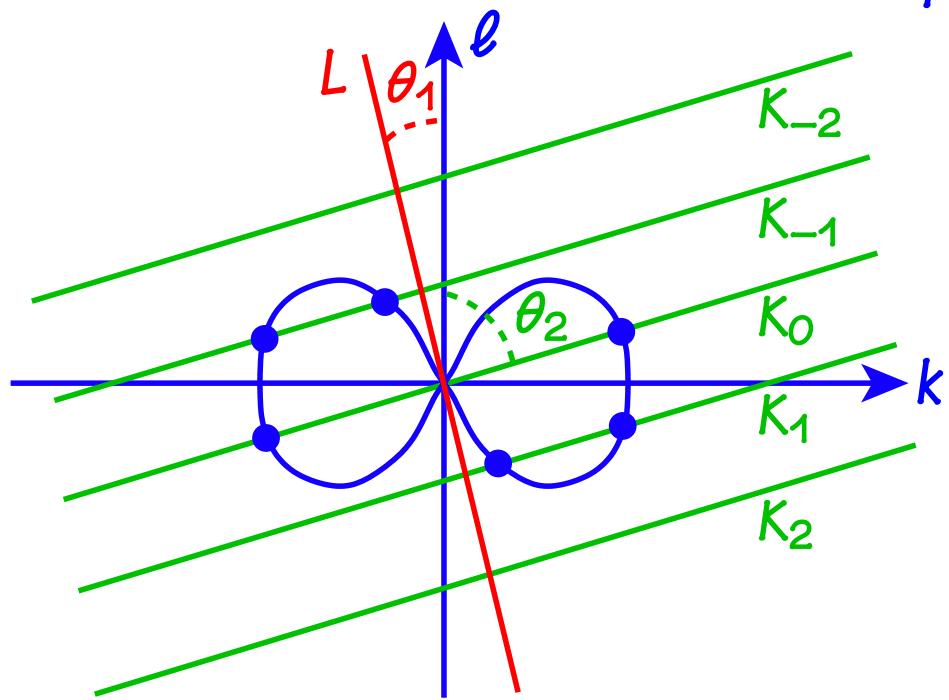
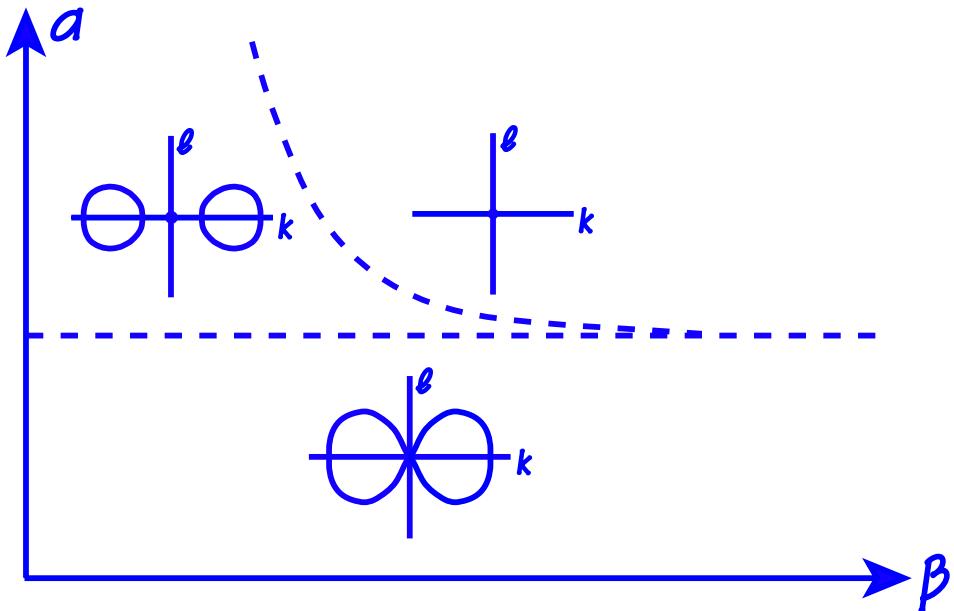
IMAGINARY EIGENVALUES

A ‘mode n ’ eigenvalue ik arises when

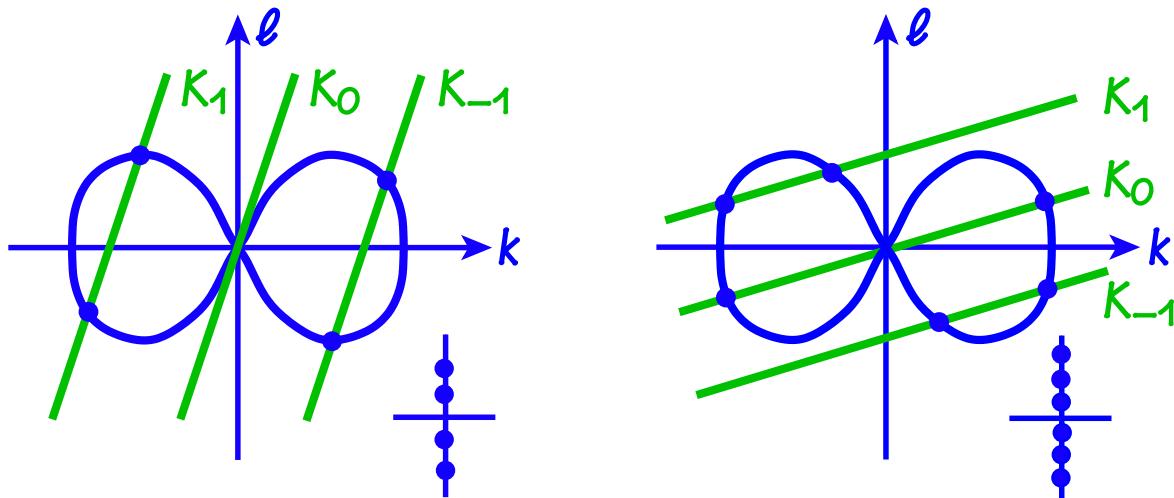
$$\begin{aligned}\eta(\xi, Z) &\sim e^{ik\xi} e^{in\nu Z} \\ &= e^{ikx} e^{ilz}, \quad k = \kappa \sin \theta_2 + n\nu \sin \theta_1 \\ &\quad l = -\kappa \cos \theta_2 - n\nu \cos \theta_1\end{aligned}\left.\right\} K_n$$

and exists when k, l satisfy the dispersion relation

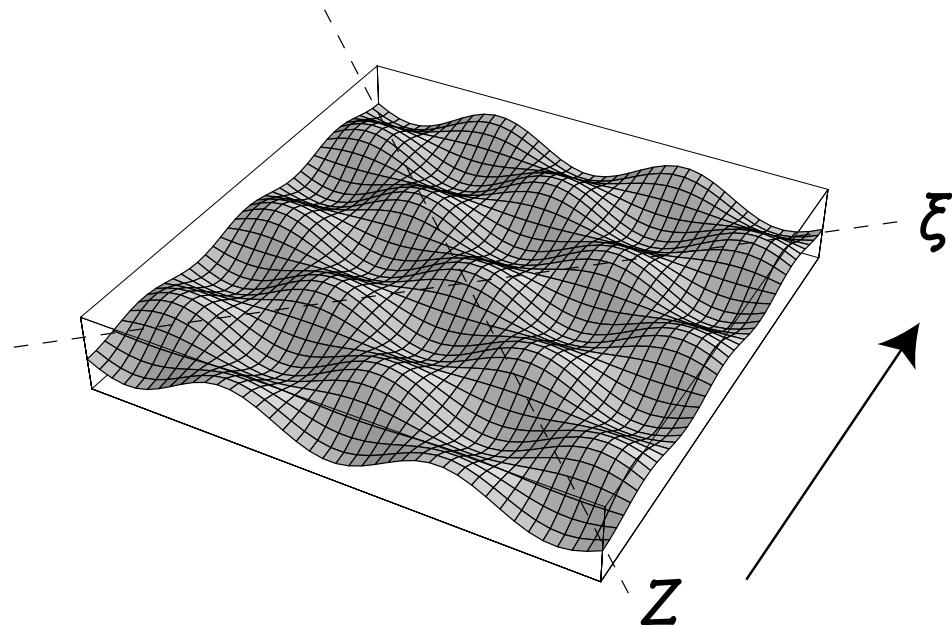
$$-k^2 + (a + \beta(k^2 + l^2))\sqrt{k^2 + l^2} \tanh \sqrt{k^2 + l^2} = 0$$



PERIODIC WAVES



- A reduced Hamiltonian system with finitely many degrees of freedom
- Periodic solutions are found using the Lyapunov centre theorem (many resonant cases)
- Inverse result (Groves & Haragus, Craig & Nicholls): Given any angles θ_1, θ_2 and frequencies κ, ν there is a corresponding doubly periodic wave

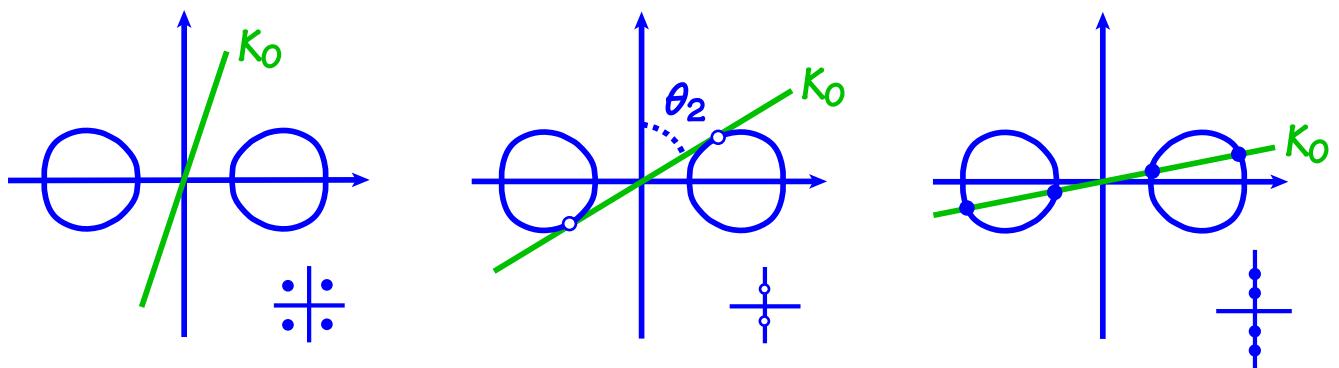




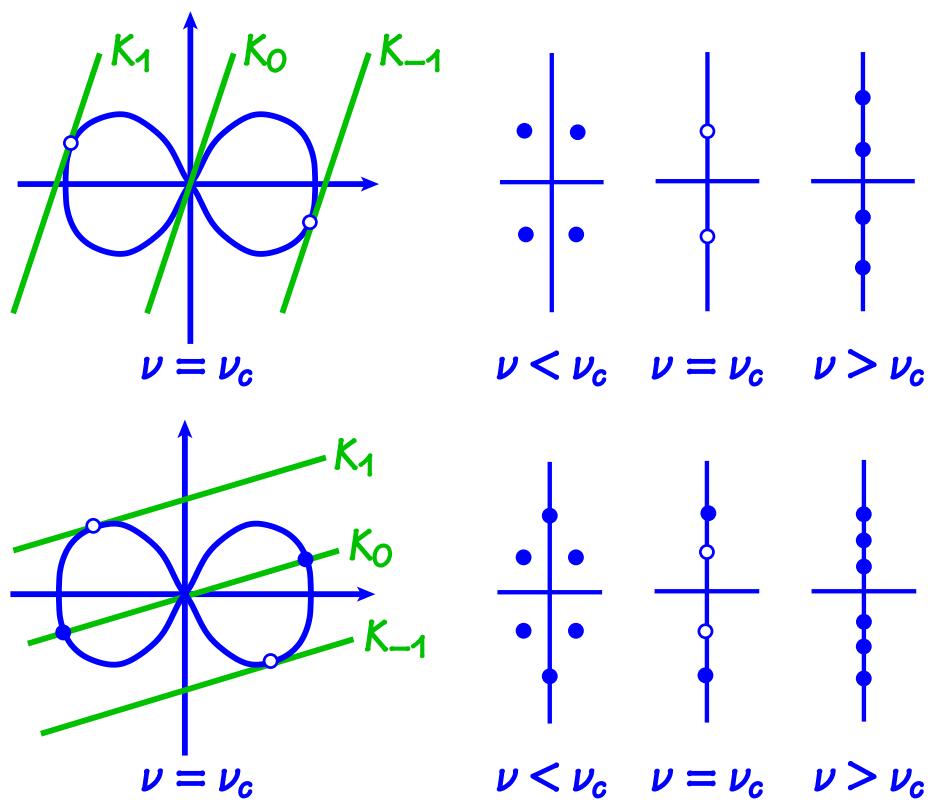


BIFURCATIONS

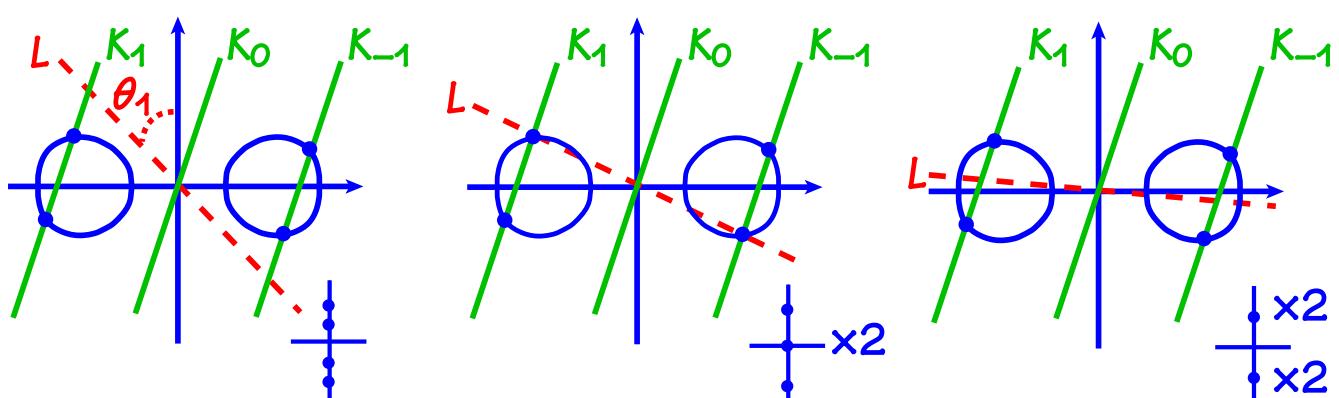
- Vary θ_2 :



- Vary ν :

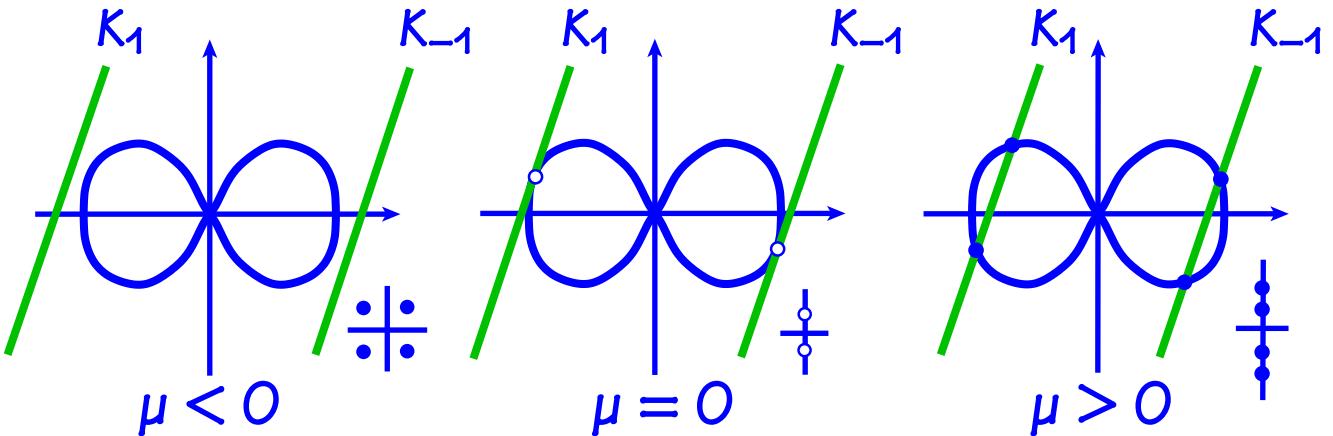


- Vary θ_1 :



- Practically every bifurcation and eigenvalue resonance possible (Groves & Haragus)

HAMILTONIAN HOPF BIFURCATION



Iooss & Kirchgässner:

- Introduce complex coordinates and use a partial normal-form transformation:

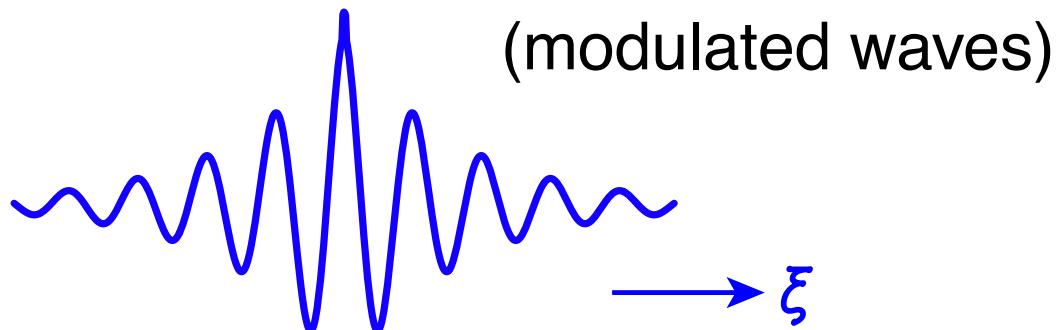
$$\dot{A} = \frac{\partial H}{\partial \bar{B}}, \quad \dot{B} = -\frac{\partial H}{\partial \bar{A}}$$

$$H = i\kappa(A\bar{B} - \bar{A}B) + |B|^2 + \underbrace{\tilde{H}(\mu, |A|^2, i(A\bar{B} - \bar{A}B))}_{\begin{array}{c} \text{Normal form} \\ \text{Polynomial of order } N+1 \end{array}} + \underbrace{O(N+2)}_{\text{Remainder term}}$$

- Without the remainder terms:
 - Complete integrability
 - A circle of homoclinic solutions for $\mu > 0$
- Two homoclinic solutions are symmetric and survive re-introduction of the remainder terms

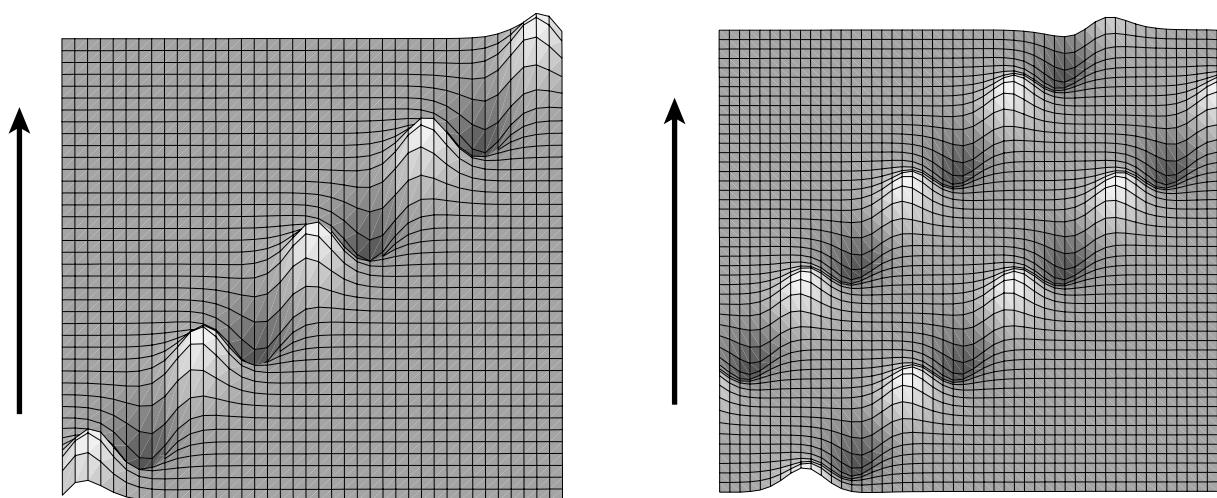
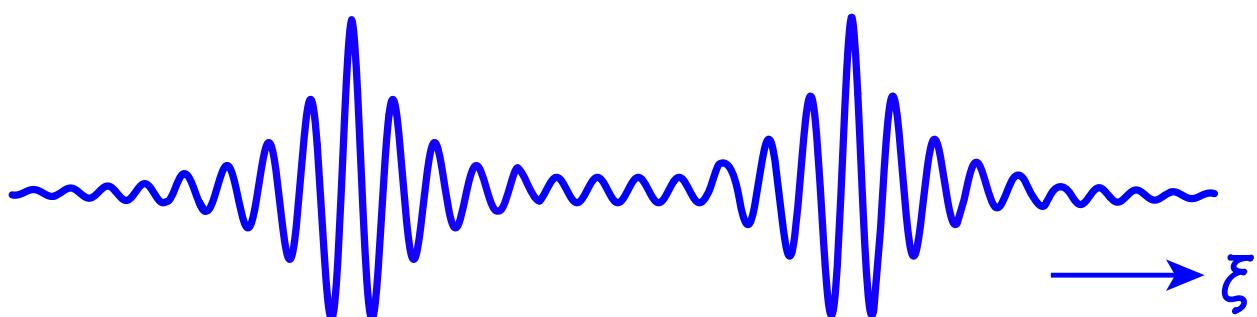
HOMOCLINIC SOLUTIONS

Iooss & Kirchgässner: A pair of symmetric homoclinic solutions



Buffoni & Groves:

- Homoclinic solutions are critical points of a functional of ‘mountain-pass’ type
- Use the calculus of variations and topological degree to construct multipulse solutions



FULLY LOCALISED SOLITARY WAVES

- Apply the direct methods of the calculus of variations to the original variational principle:

$$\delta \int_{\mathbb{R}^2} \left\{ \int_0^1 \left(\frac{1}{2} \left(\Phi_x - \frac{y\eta_x \Phi_y}{1+\eta} \right)^2 + \frac{\Phi_y^2}{2(1+\eta)^2} + \frac{1}{2} \left(\Phi_z - \frac{y\eta_z \Phi_y}{1+\eta} \right)^2 \right) (1+\eta) dy + \eta \Phi_x|_{y=1} + \frac{1}{2} a\eta^2 + \beta [\sqrt{1+\eta_x^2 + \eta_z^2} - 1] \right\} dx dz = 0$$

- Quasilinear structure
- Groves & Sun: Reduce to a locally equivalent semilinear problem using a generalisation of the variational Lyapunov-Schmidt reduction procedure

